

# Determination of hardness to elastic modulus ratios using Knoop indentation measurements and a model based on loading and reloading half-cycles

J. C. CONWAY, Jr

*Department of Engineering Science and Mechanics, The Pennsylvania State University, University Park, Pennsylvania 16802, USA and Ceramic Finishing Co., State College, Pennsylvania 16804, USA*

A model is developed which results in an expression relating the residual in-surface dimension of the minor diagonal of a Knoop indenter through a "residual-width parameter"  $(b_R/b^*)^2$  to the hardness-to-modulus ratio  $(H/E)$  of elastic-plastic materials. The relationship is shown to predict within reasonable accuracy both the intercept and slope of a plot of  $(b_R/b^*)^2$  against  $(H/E)$ , using data for a variety of materials. The work supplements concepts presented in two previous investigations.

## 1. Introduction

Measurement of the hardness-to-modulus ratio  $H/E$  of elastic-plastic materials is an important prerequisite to the determination of fracture toughness by indentation methods [1]. Two indentation techniques have been developed to date to measure this important parameter. The first, developed by Lawn and Howes [2], demonstrated that  $H/E$  could be related to the elastic depth recovery of Vickers indentations. The model adopted was based on maintaining compatibility at the maximum penetration depth for an elastic-plastic loading half-cycle and an elastic unloading or reloading half-cycle. The second technique, developed by Marshall *et al.* [3], related  $H/E$  to the residual in-surface dimension of the minor diagonal of a Knoop indentation. This method was based on a model utilizing the superposition of solutions for an elliptical hole subjected to uniaxial stress and the fact that the in-surface dimension of the major diagonal undergoes relatively little length change during unloading. Both techniques were shown to give accurate predictions of  $H/E$  through the results of indentation experiments conducted on selected materials, mainly ceramics.

This paper applies the methodology of the Lawn and Howes model [2] to a Knoop indentation and offers a supplemental approach to that adopted by Marshall *et al.* [3]. The model is based on an elastic-plastic loading half-cycle and an elastic unloading or reloading half-cycle, and  $H/E$  is related to the residual in-surface dimension of the minor diagonal of a Knoop indentation.

## 2. Model

Referring to Fig. 1, during the loading half-cycle the material immediately surrounding the indenter undergoes elastic-plastic deformation resulting in the

characteristic in-surface contact dimensions  $2a^*$  and  $2b^*$  at maximum indentation load,  $P^*$ . As unloading occurs, elastic recovery takes place along the in-surface dimension of the minor diagonal resulting in a residual dimension  $2b_R$ . Relatively little change takes place in the in-surface dimension of the major diagonal during unloading.

During indentation, the projected area is  $A_p = 2ab$  and the mean contact pressure is  $p_0 = P/2ab$  with  $P$  being the contact load. Adopting the notation of Lawn and Howes [2] with  $a = 7.11b$  and  $p_0 = H$  during the loading half-cycle,

$$P = 14.22Hb_H^2 \quad (1)$$

In deriving this expression, "pile-up" or "sink-in" of the plastically-deformed material surrounding the contact zone has been neglected. As can be seen from Equation 1, the indentation mechanics is controlled by the hardness  $H$  during the elastic-plastic loading half-cycle.

During the elastic unloading or reloading half-cycle, the mean contact pressure associated with a rigid wedge or cone is adopted as  $p_m = E \times \cot \gamma/2(1 - \nu^2)$  [4]. In this expression,  $E$  and  $\nu$  are the modulus of elasticity and Poisson's ratio of the indented material, respectively, and  $\gamma$  is the average half-angle of a Knoop indenter  $(\alpha + \beta)/2$ . Elastic loading of a half-space may then be generally described as

$$P_E = \left[ \left( \frac{14.22E}{2(1 - \nu^2)} \right) \cot \gamma \right] b_E^2 \quad (2)$$

where  $2b_E$  is the in-surface dimension of the minor diagonal at any given load  $P_E$ .

Referring to Fig. 2, the loading half-cycle is shown as OC and is described by Equation 1. The unloading half-cycle is shown as CD, resulting in a residual

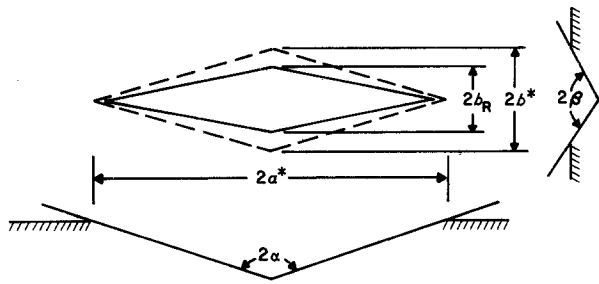


Figure 1 Knoop indentation impression showing in-surface elastic recovery of minor diagonal.

in-surface minor diagonal dimension  $2b_R$ . Since the unloading and reloading stages are reversible, reloading along path  $D'C'$  will be used to describe the unloading half-cycle. Reloading can be described by writing Equation 2 referenced with respect to a coordinate system with origin at  $D'$  as

$$P' = \left[ \left( \frac{14.22E}{2(1-\nu^2)} \right) \cot \gamma \right] (b_E^2 - b_R^2) \quad (3)$$

where  $b_R \leq b_E \leq b^*$  during reloading and  $2b^*$  is the in-surface dimension of the minor diagonal at maximum load  $P^*$ . This is equivalent to adopting the "compressed spring" model of Lawn and Howes [2]. Compatibility conditions at the end of the loading and reloading half-cycles require that at  $P = P^*$ ,  $b_H = b_E = b^*$  so that Equations 1 and 3 give

$$\left( \frac{b_R}{b^*} \right)^2 = 1 - [2(1 - \nu^2) \tan \gamma] \left( \frac{H}{E} \right) \quad (4)$$

This expression can, of course, be written in terms of the in-surface major diagonal dimension at maximum load  $2a^*$  since  $a^* = 7.11b^*$ . In conducting actual tests, this may be the more convenient formulation since the length of the in-surface major diagonal undergoes relatively little dimensional change on unloading. In either case, Equation 4 indicates that  $H/E$  can be directly related to the residual in-surface minor diagonal dimension  $2b_R$ .

### 3. Results and discussion

The relationship between the hardness-to-modulus ratio  $H/E$  and the residual in-surface minor diagonal dimension  $2b_R$  as expressed in Equation 4 was verified by plotting existing data for a variety of materials as  $(b_R/b^*)^2$  against  $H/E$  in Fig. 3. The data were obtained from Table I and Fig. 2 of Marshall

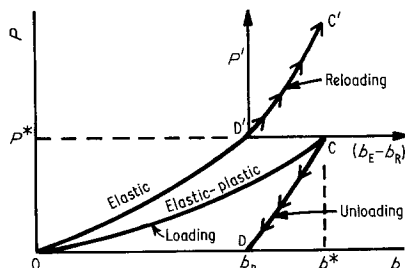


Figure 2 Loading, unloading and reloading half-cycles.  $P$  is the contact load and  $b$  the minor diagonal in-surface half dimension. Loading plot from Equation 1; reloading plot from Equation 2.

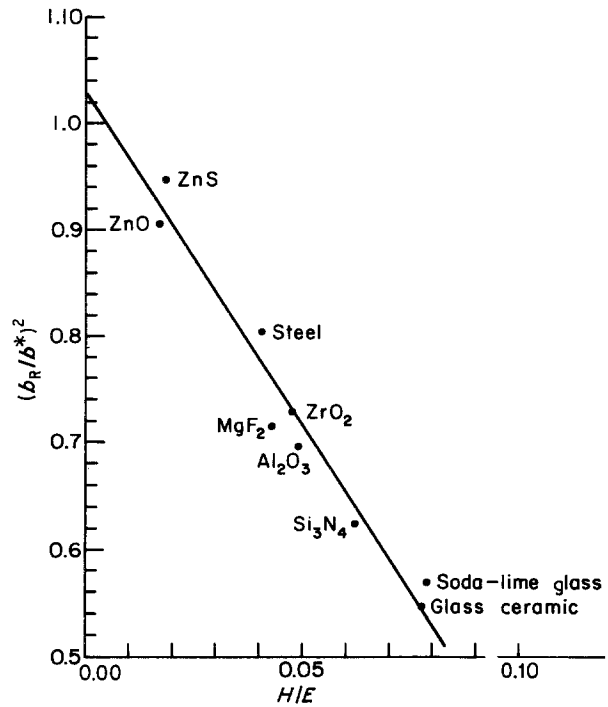


Figure 3 Residual width parameter  $(b_R/b^*)^2$  as a function of hardness-to-modulus ratio ( $H/E$ ).

*et al.* [3] and are shown in Table I of this paper. Values of  $(b_R/a')$  graphically shown by Marshall *et al.* [3] were converted to the "residual-width parameter"  $(b_R/b^*)^2$  by assuming that  $a' \cong a^* = 7.11b^*$ . The solid line through the data is a least-squares curve fit with an intercept of 1.03 and a slope of  $-6.19$ . The intercept and slope predicted from Equation 4, assuming  $\nu = 0.25$  and  $\gamma = 75^\circ$  (the average half-angle of a Knoop indenter) are 1.00 and  $-6.99$  respectively. These are reasonable results when considering the sensitivity of the quantity  $2(1 - \nu^2) \tan \gamma$  to changes in the average half-angle of the indenter  $\gamma$ , as well as the data scatter. The data do not span the elastic-plastic spectrum but ideal plastic behaviour is indicated in Fig. 3 at the point  $(b_R/b^*)^2 = 1.00$ ,  $H/E = 0$ . Extrapolation of the curve fit to ideal elastic behaviour gives  $(b_R/b^*)^2 = 0$ ,  $H/E = 0.17$ .

Results obtained by using the model presented in this paper compare well with those obtained by Marshall *et al.* [3], as can be seen by comparing like intercepts and slopes of plots for a residual-width parameter against hardness-to-modulus ratio for like data. Converting the ordinate of Fig. 3 of this work from  $(b_R/b^*)^2$  to  $b_R/a^* = b'/a'$  for  $b^* = 7.11a^*$  gives intercept of 0.14 and a slope of  $-0.48$ , which compare

TABLE I Material and dimensional parameters

Material	$H/E$	$b_R/a^*$	$(b_R/b^*)^2$
Soda-lime glass	0.079	0.106	0.568
Glass-ceramic	0.078	0.104	0.547
$Si_3N_4$ (hot-pressed)	0.062	0.111	0.623
$Al_2O_3$ (hot-pressed)	0.050	0.118	0.704
$ZrO_2$ (partially stabilized)	0.048	0.120	0.728
$MgF_2$ (hot-pressed)	0.043	0.119	0.716
Steel (hardened)	0.041	0.126	0.803
ZnS (hot-pressed)	0.019	0.137	0.949
ZnO (hot-pressed)	0.017	0.134	0.908

favourably with values of 0.14 and  $-0.45$  obtained by Marshall *et al.* [3]. The model described in this paper, however, avoids the assumption of a two-dimensional elliptical hole solution and allows direct analytical determination of the slope and intercept with reasonable accuracy.

### Acknowledgements

The author thanks H. P. Kirchner of Ceramic Finishing Co. for many helpful discussions leading to this work and for his comments on the manuscript. The work was sponsored by the US Department of Energy under Contract No. DE-AC01-83ER80015.

### References

1. G. R. ANSTIS, P. CHANTIKUL, B. R. LAWN and D. B. MARSHALL, *J. Amer. Ceram. Soc.* **64** (1981) 533.
2. B. R. LAWN and V. R. HOWES, *J. Mater. Sci.* **16** (1981) 2745.
3. D. B. MARSHALL, T. NOMA and A. G. EVANS, *Commun. Amer. Ceram. Soc.* **65** (1982) C175.
4. K. L. JOHNSON, *J. Mech. Phys. Solids* **18** (1970) 118.

*Received 29 July*

*and accepted 18 September 1985*